

## Numerical methods for hyperbolic systems

### Exercise sheet 2: Linear hyperbolic systems

#### Exercise 1

We consider the wave equation

$$\begin{cases} \partial_t p + \partial_x u = 0, \\ \partial_t u + \partial_x p = 0, \end{cases} \quad (1)$$

1. Diagonalize the system and show that the upwind scheme is given by

$$\begin{cases} \frac{v_j^{n+1} - v_j^n}{\Delta t} + \frac{v_j^n - v_{j-1}^n}{\Delta x} = 0, \\ \frac{w_j^{n+1} - w_j^n}{\Delta t} - \frac{w_{j+1}^n - w_j^n}{\Delta x} = 0, \end{cases} \quad (2)$$

with  $v = p + u$  and  $w = p - u$ .

2. Prove that the upwind scheme (2) for the initial system (1) can be written on the following form

$$\begin{cases} \frac{p_j^{n+1} - p_j^n}{\Delta t} + \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} - \frac{p_{j+1}^n - 2p_j^n + p_{j-1}^n}{2\Delta x} = 0, \\ \frac{u_j^{n+1} - u_j^n}{\Delta t} + \frac{p_{j+1}^n - p_{j-1}^n}{2\Delta x} - \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{2\Delta x} = 0, \end{cases} \quad (3)$$

and used this form to compute the consistency error.

3. Prove that the scheme (2) satisfies the maximum principle under a CFL condition.

4. Prove that the scheme (2) is stable for all  $l^q$  norms ( $1 \leq q \leq \infty$ ) using the previous result and convex functions. The  $l^q$  norm is defined by

$$\|(v, w)\|_{l^q} = \left( \Delta x \sum_j \|(v_j, w_j)\|_q^q \right)^{\frac{1}{q}}$$

with  $\|(v_j, w_j)\|_q^q = |v_j|^q + |w_j|^q$ .

**Additional question**

We introduce the damped wave equation

$$\begin{cases} \partial_t p + \partial_x u = 0, \\ \partial_t u + \partial_x p = -\sigma u, \end{cases} \quad (4)$$

and the upwind scheme associated

$$\begin{cases} \frac{p_j^{n+1} - p_j^n}{\Delta t} + \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} - \frac{p_{j+1}^n - 2p_j^n + p_{j-1}^n}{2\Delta x} = 0, \\ \frac{u_j^{n+1} - u_j^n}{\Delta t} + \frac{p_{j+1}^n - p_{j-1}^n}{2\Delta x} - \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{2\Delta x} = -\sigma u_j^n, \end{cases} \quad (5)$$

5. We call "steady states" the solutions of the systems defined by  $\partial_x u = 0$  and  $\partial_x p = -\sigma u$ . Prove that (5) preserve exactly the steady states.

## Exercise 2

We consider the Maxwell equation

$$\begin{cases} \mu \partial_t B + c \partial_x E = -c \sigma^* B, \\ \varepsilon \partial_t E + c \partial_x B = -c \sigma E, \end{cases} \quad (6)$$

with periodic boundary condition on  $\Omega = [0, L]$  and  $\mu > 0$ ,  $\varepsilon > 0$ ,  $\sigma > 0$ ,  $\sigma^* > 0$ ,  $c > 0$ .

1. Prove the following energy equality and the uniqueness of the solutions.

$$\frac{d}{dt} \left( \int_{\Omega} \varepsilon |E(t, x)|^2 + \mu |B(t, x)|^2 dx \right) = - \int_{\Omega} \sigma |E(t, x)|^2 + \sigma^* \mu |B(t, x)|^2 dx \quad (7)$$

2. We introduce the plane waves (which are a good approximations of physical waves) defined by  $E(t, x) = E_0 e^{j(wt - kx)}$  and  $B(t, x) = B_0 e^{j(wt - kx)}$  ( $j$  complex number) with  $E_0 \in \mathbb{R}$ ,  $B_0 \in \mathbb{R}$ ,  $k$  the wave vector and  $w$  the frequency. Give the conditions (called dispersion relation) on  $w$  and  $k$  such as the plane waves are solutions of (6) for  $\sigma = 0$  and  $\sigma^* = 0$ .

In this part we assume that  $\sigma = 0$  and  $\sigma^* = 0$ . Now we introduce the DG centered scheme for (6). The mesh  $\Omega_h$  is defined by  $n + 1$  points  $x_i$  and  $n$  cells  $K_i = [x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}]$ . The volume of the cell  $K_i$  is  $\Delta x_i = |x_{i+\frac{1}{2}} - x_{i-\frac{1}{2}}|$ . We call a generic cell  $K$ . To finish the test function are defined by  $v \in V_h = \{v/v|_K \in \mathbb{P}^p(K)\}$ . The scheme is given by

$$\begin{cases} \varepsilon \sum_{l=0}^k \int_{K_i} \phi_l^i \phi_m^i \left( \frac{E_{l,i}^{n+1} - E_{l,i}^n}{\Delta t} \right) - c \sum_{l=0}^k B_{l,i}^n \int_{K_i} \phi_l^i \partial_x \phi_m^i + c \sum_{l=0}^k [B \phi_m^i]_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} = 0, & \forall 0 \leq m \leq k, \\ \mu \sum_{l=0}^k \int_{K_i} \phi_l^i \phi_m^i \left( \frac{B_{l,i}^{n+1} - B_{l,i}^n}{\Delta t} \right) - c \sum_{l=0}^k E_{l,i}^n \int_{K_i} \phi_l^i \partial_x \phi_m^i + c \sum_{l=0}^k [E \phi_m^i]_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} = 0, & \forall 0 \leq m \leq k, \end{cases} \quad (8)$$

with  $[B \phi_m^i]_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} = \frac{1}{2} \left( B_{l,i+1}^n \phi_l^{i+1}(x_{i+\frac{1}{2}}) \phi_m^i(x_{i+\frac{1}{2}}) + B_{l,i}^n \phi_l^i(x_{i+\frac{1}{2}}) \phi_m^i(x_{i+\frac{1}{2}}) \right)$

$$-\frac{1}{2} \left( B_{l,i-1}^n \phi_l^{i-1}(x_{i-\frac{1}{2}}) \phi_m^i(x_{i-\frac{1}{2}}) + B_{l,i}^n \phi_l^i(x_{i-\frac{1}{2}}) \phi_m^i(x_{i-\frac{1}{2}}) \right).$$

**3.** We consider  $V_h = P^1(K)$ . We propose to use the Lagrange polynomial associated with the point  $x_{i-\frac{1}{2}}$  and  $x_{i+\frac{1}{2}}$ . Prove that the family is a basis of  $V_h$ . Write the scheme in a cell  $K_i$ .

**4.** In this exercise we propose to study the numerical dispersion relation which define the numerical wave vector  $\tilde{k}$  for  $V_h = P^0(K) = \text{Span}(1)$ . Write the scheme for this basis. Now we define  $B_i^n = B_0 e^{j(wn\Delta t - \tilde{k}i\Delta x)}$  and  $E_i^n = E_0 e^{j(wn\Delta t - \tilde{k}i\Delta x)}$  with  $j$  the complex number and  $i$  the index of the cell. Gives the relation between  $w$  and  $\tilde{k}$  such as the discrete plane waves are solutions of (8). Show that the numerical dispersive relation is  $\tilde{k}^2 = \frac{w^2}{c^2} + O(\Delta x^p + \Delta t^q)$  with  $p > 1$  and  $q > 1$ .